

# Reduced Discrete Variational Multi-body Mechanics

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## Abstract

This paper derives a numerical method for simulating multi-body mechanics using reduced discrete variational principles. Motivated by the favorable nature of symplectic integrators for multi-body systems [1, 2, 3, 4] we construct, for the first time, integrators defined through reduced variational principles. Then we employ the framework to compute motion controls using discrete optimal control in shape-space to satisfy boundary conditions of the full system.

*Tree/Graph Structure.* Consider a mechanical system consisting of interconnected rigid bodies indexed using the set of integers so that joint  $\#j$  connect body  $\#a_j$  to body  $\#b_j$ . Without loss of generality assume that body  $i = 0$  is the reference body, i.e. it has no incoming joints.

*Rigid Body Configuration Space.* The configuration of body  $\#i$  is denoted by  $g_i \in SE(3)$ , while its body-fixed velocity is denoted  $\xi_i \in \mathfrak{se}(3)$  and defined by  $\xi_i = g_i^{-1} \dot{g}_i$ . In general, an element  $g \in SE(3)$  and its inverse are defined using a matrix  $R \in SO(3)$  and vector  $x \in \mathbb{R}^3$  according to

$$g = \begin{pmatrix} R & x \\ 0 & 1 \end{pmatrix}, \quad g^{-1} = \begin{pmatrix} R^T & -R^T x \\ 0 & 1 \end{pmatrix}.$$

The velocities  $\xi \in \mathfrak{se}(3)$  are identified with body-fixed angular and linear velocities denoted  $\omega \in \mathbb{R}^3$  and  $v \in \mathbb{R}^3$ , respectively, through

$$\xi = \begin{pmatrix} \widehat{\omega} & v \\ 0 & 0 \end{pmatrix}, \quad \text{where } \widehat{\omega} = \begin{bmatrix} 0 & -\omega^3 & \omega^2 \\ \omega^3 & 0 & -\omega^1 \\ -\omega^2 & \omega^1 & 0 \end{bmatrix}.$$

The algorithm is thus implemented in terms of vectors in  $\mathbb{R}^6$  rather than matrices in  $\mathfrak{se}(3)$ . In the sequel we will use the following shorthand notation  $g_{0:N} = \{g_0, \dots, g_N\}$  to denote the configurations of all bodies of the multi-body system.

The inertia tensor of the body is denoted by  $\mathbb{J}_i : \mathfrak{se}(3) \rightarrow \mathfrak{se}(3)^*$  and implemented as a 6x6 diagonal matrix. Each body is subject to potential energy, e.g. due to gravity, defined by the function  $V : SE(3) \rightarrow \mathbb{R}$ .

*Joints.* The system has  $n$  joints described by parameters  $r \in M$ , where  $M \subset \mathbb{R}^n$  is the *shape space*. The relative transformation between the reference body configuration  $g$  and the configuration of body  $\#i$  is denoted by  $g_{0i} : M \rightarrow SE(3)$ . We assume that all joints are controlled using control inputs denoted by  $f \in \mathbb{R}^m$ .

*Discrete Trajectory.* The trajectory of a rigid body over the time interval  $[0, T]$  is represented numerically using a set of  $K + 1$  equally spaced in time points denoted  $g_{0:K} := \{g_0, \dots, g_K\}$ , where  $g_k \approx g(kh)$  and  $h = T/K$  denotes the time-step. The curve segment between each pair of points  $g_k$  and  $g_{k+1}$  is interpolated by a short curve that must lie on the manifold  $SE(3)$ . The simplest way to construct such a curve is through the map  $\tau$  and a vector  $\xi_k \in \mathfrak{se}(3)$  such that  $\xi_k = \tau^{-1}((g_k)^{-1}g_{k+1})/h$ . Such a discretization can be regarded as a discrete approximation of the continuous curve according to

$$g(t) \approx g_k \tau((t - hk)\xi_k), \quad r(t) \approx r_k + \frac{t - hk}{h} (r_{k+1} - r_k), \quad \text{for } t \in [kh, (k+1)h].$$

*Reduced Lagrangian.* The dynamics will be derived using the reduced Lagrangian [5, 6, 4]  $\ell : G \times \mathfrak{g} \times TM \rightarrow \mathbb{R}$  defined by

$$\ell(g, \Omega, r, \dot{r}) = \frac{1}{2} \begin{bmatrix} \Omega \\ \dot{r} \end{bmatrix}^T \begin{bmatrix} \mathbb{I}(r) & 0 \\ 0 & m(r) - \mathcal{A}(r)^T \mathbb{I}(r) \mathcal{A}(r) \end{bmatrix} \begin{bmatrix} \Omega \\ \dot{r} \end{bmatrix} - V(g, r) \quad (1)$$

where  $\mathcal{A} : TQ \rightarrow \mathfrak{g}$  is the mechanical connection defined by  $g^{-1}\dot{g} = \Omega - \mathcal{A}(r)\dot{r}$  and  $\mathbb{I} : \mathfrak{g} \rightarrow \mathfrak{g}^*$  is the locked inertia tensor. Regarding them as matrices they are computed according to

$$\mathbb{I} = \mathbb{I}_0 + \sum_{i=1}^n A_i^T \mathbb{I}_i A_i, \quad m = \sum_{i=1}^n J_i^T \mathbb{I}_i J_i, \quad \mathcal{A} = \mathbb{I}^{-1} \left( \sum_{i=1}^n A_i^T \mathbb{I}_i J_i \right), \quad (2)$$

using the adjoint notation  $A_i := \text{Ad}_{g_{0i}^{-1}(r)}$ , and Jacobian  $J_i := g_{0i}^{-1}(r) \partial_r g_{0i}(r)$ .

*Discrete Variational Principle.* With a discrete trajectory and a Lagrangian in place we construct a structure-preserving (i.e. group, momentum, and symplectic) integrator of the dynamics. This is accomplished by defining a discrete variational principle through a trapezoidal quadrature approximation and, using the shorthand

$$\ell_{k+\alpha} := (1 - \alpha)\ell(g_k, \Omega_k, r_k, u_k) + \alpha\ell(g_{k+1}, \Omega_k, r_{k+1}, u_k), \quad \mathcal{A}_{k+\alpha} := (1 - \alpha)\mathcal{A}(r_k) + \alpha\mathcal{A}(r_{k+1}), \quad (3)$$

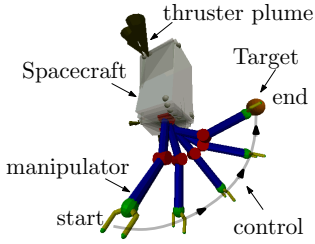
results in the discrete equations of motion (see [4])

$$(d\tau_{\Omega_k - \mathcal{A}_{k+\alpha} u_k}^{-1})^* \mu_k - (d\tau_{-\Omega_{k-1} + \mathcal{A}_{k-1+\alpha} u_{k-1}}^{-1})^* \mu_{k-1} = -h g_k^* \partial_g V(g_k, r_k), \quad (4)$$

$$(p_k - p_{k-1})/h - (1 - \alpha) \langle \mu_k, \partial \mathcal{A}(r_k)(\cdot, u_k) \rangle - \alpha \langle \mu_{k-1}, \partial \mathcal{A}(r_k)(\cdot, u_{k-1}) \rangle = ((1 - \alpha) \partial_r \ell_k^- + \alpha \partial_r \ell_{k-1}^+) + f_k, \quad (5)$$

where  $\Omega_k := \mathbb{I}^{-1} \mu_k$  and  $u_k := \partial_u \ell_{k+\alpha}^*(p_k - \mathcal{A}_{k+\alpha}^* \mu_k)$ . The discrete Euler-Poincare equation (4) and the discrete reduced Euler-Lagrange equation (5) are implicit equations in terms of the unknown momenta  $\mu_k$  and  $p_k$ . In addition, once (4)-(5) are solved the updated configurations are computed according to

$$r_{k+1} = r_k + h u_k, \quad g_{k+1} = g_k \tau(h(\Omega_k + ((1 - \alpha)\mathcal{A}(r_k) + \alpha\mathcal{A}(r_{k+1})) u_k)). \quad (6)$$



*Computed Example.* Consider a robotic satellite equipped with a 3-dof manipulator with two joints: 2-dof shoulder and 1-dof elbow. The general task is to compute a minimum control effort trajectory reaching a desired state. The algorithm is implemented and simulated as shown on left.

*Motion Control.* The reduced formulation enables a natural approach to motion control. Imagine that the free-floating manipulator must grasp a particular target in task-space, or equivalently to move between two given zero-velocity states  $(g_0, 0, r_0, 0)$  and  $(g_N, 0, r_N, 0)$ . The approximate solution shown above was computed by performing nonlinear optimization of shape trajectories  $r_{0:K}$  and by reconstructing the full state  $(g_{k+1}, \xi_k, r_{k+1}, u_k)$  for each  $k$  using (6). The objective function includes the control effort as well as a final cost penalizing deviation from the given boundary conditions.

Motion optimization based on the mechanical connection is efficiently implemented using a stage-wise Newton method [7] which has complexity linear in the number of discrete segments  $K$ . The algorithm requires first-order derivatives of the equations of motion (4)-(6) and first and second-order derivatives of the given cost function.

## References

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